

1 Frequency Domain EOF  $\rightarrow$  phase propagation of waves/oscillations  
(time domain)  $\rightarrow$  time lag problems

conventional EOF detects standing oscillations only  
"patterns" are either in or out of phase.

complex (frequency domain) EOF detects propagating and standing oscillations  
"patterns" can move, i.e., time-space lagged correlations

I. brute force approach on <sup>set of</sup> series  $\{x\} = \{x_1(t), x_2(t), \dots, x_N(t)\}$   
(complex)

1. set up cross-spectral matrix

set-up covariance matrix

$$S_{ij}(f) = \frac{1}{T} E[X_i^*(f) X_j(f)]$$

$$R_{ij} = \langle x_i(t) x_j(t) \rangle$$

2. find eigenvalues of matrix  $S_{ij}(f)$

find eigenvalues of matrix  $R_{ij}$

$$\rightarrow \lambda_i = \lambda_i(f) \quad i=1,2,\dots,N$$

$$\rightarrow \lambda_i \quad i=1,2,\dots,N$$

(power spectra)

(variance explained)

as  $S_{ij}$  is Hermitian, ~~the eigen~~  
the eigen values are real

as  $R_{ij}$  is symmetric  
the eigen values are real

3. find eigenvectors of matrix  $S_{ij}$   
for each eigenvalue  $\lambda_i(f)$

find eigenvectors of matrix  $R_{ij}$   
for each eigenvalue  $\lambda_i$

eigenvectors are now complex  
different set of vectors at each frequency

4. find the amplitude series  
for each eigenvector

find the amplitude series

$$A_n(f_2) = \sum_{i=1}^M X_{ij}(x_i, f_2) \cdot \phi_n(x_i)$$

$$A_n(t_2) = \sum_{i=1}^M \alpha_i(x_i, t_2) \phi_n(x_i)$$

Michelson (1982)  
JPO 694-703



Brute force approach to frequency domain EOF

A set of observations are

$$\vec{u}(t) = \{u_1(t, x_1), u_2(t, x_2), \dots, u_M(t, x_M)\} = \{u_1(t), u_2(t), \dots, u_M(t)\}$$

An EOF analysis looks at

$$R_{pq}(\tau=0) = \langle u_p(t) u_q(t) \rangle$$

while FEOF analysis looks at

$$\tilde{R}_{pq}(\tau) = \langle u_p(t) u_q(t+\tau) \rangle$$

in the form

$$\tilde{S}_{pq}(f) = \int_{-\infty}^{\infty} \tilde{R}_{pq}(\tau) e^{-j2\pi f\tau} d\tau \equiv \frac{1}{T} U_p^*(f) U_q(f)$$

The EOF (FEOF) analysis tries to find the real (real) eigenvalues  $\lambda_i$  and corresponding real (complex) eigenvectors  $\vec{\phi}_i$  of the real (complex) and symmetric (Hermitian) matrix  $R_{pq} = R_{qp}$  ( $\tilde{S}_{pq} = \tilde{S}_{qp}^*$ ), i.e.,

$$R_{pq} \cdot \vec{\phi}_i = \lambda_i \vec{\phi}_i \quad \left( \tilde{S}_{pq} \cdot \vec{\phi}_i = \lambda_i \vec{\phi}_i \right) \quad i=1,2,\dots,M$$

The <sup>real</sup> EOF amplitude functions are  $A_n(t) = \sum_{i=1}^M u_i(t) \phi_n(x_i)$

while the <sub>complex</sub> FEOF amplitudes are

$$A_n(t) = \sum_{i=1}^M u_i(t) \int_{-\infty}^{\infty} \phi_n(x_i, f) e^{j2\pi f t} df$$



II. "Wide Band" frequency domain EOF (Barnett (1983), Monthly Weather Review p. 756-773 also Harel (1984), J. Climate Appl. Meteor. 1660-1678)

- widely used, especially by meteorologists and climatologists
- essentially an average over all frequencies of the cross-spectral matrix is analyzed; average can also be over a "wide" frequency band

A set of scalars (can be generalized to vectors  $\rightarrow$  Barnett (1983)) are

$$\vec{u}(t) = \{u_1(t, x_1), u_2(t, x_2), \dots, u_M(t, x_M)\} = \{u_1(t), u_2(t), \dots, u_M(t)\}$$

and have a "harmonic" or "Fourier" representation

$p$  - "location"  
 $q$  for frequency

$$u_p(t) = \sum_{q=1}^N a_p \cos(2\pi \frac{q}{T} t) + b_p \sin(2\pi \frac{q}{T} t)$$

$$= \sum_{\omega} a_p(\omega) \cos \omega t + b_p(\omega) \sin(\omega t) \quad \omega = \frac{2\pi}{T} i$$

In order to describe traveling (time lagged) feature we need a

complex representation such as

$$U_p(t) = \sum_{\omega} c_p(\omega) e^{-j\omega t} \quad c_p(\omega) = a_p(\omega) + i b(\omega)$$

$$= \sum_{\omega} [a_p(\omega) + j b_p(\omega)] [\cos \omega t - j \sin \omega t]$$

$$= \sum_{\omega} a_p(\omega) \cos \omega t + b_p(\omega) \sin \omega t + j (b_p(\omega) \cos \omega t - a_p(\omega) \sin \omega t)$$



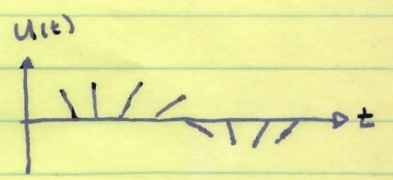
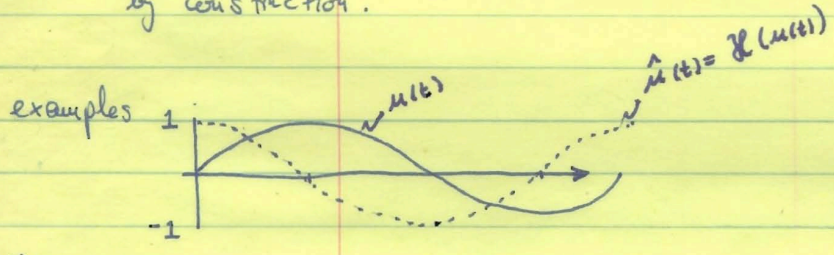
$$U_p(t) = u_p(t) + j \hat{u}_p(t)$$

$$\text{Re}(U_p) = \sum_{\omega} a_p(\omega) \cos \omega t + b_p(\omega) \sin \omega t$$

$$\begin{aligned} \text{Im}(U_p) &= \sum_{\omega} b_p(\omega) \cos \omega t - a_p(\omega) \sin \omega t \\ &= \sum_{\omega} b_p(\omega) \sin(\omega t + \pi/2) + a_p(\omega) \cos(\omega t + \pi/2) \end{aligned}$$

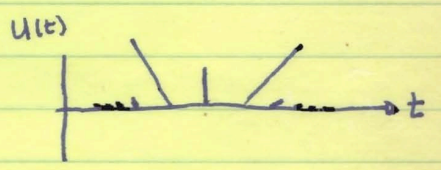
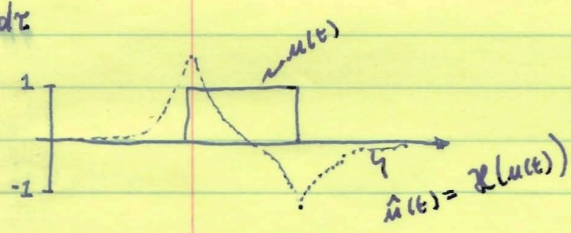
→  $\hat{u}_p(t)$  represents  $u_p(t)$  phase shifted by  $\pi/2$ ; same Fourier coefficients!

This represents formally a Hilbert transform, i.e.,  $u_p$  and  $\hat{u}_p$  are Hilbert transform pairs. Much is known about them (similar to Fourier transforms) and their characteristics. They are orthogonal by construction.



$$\hat{u}(t) = \mathcal{H}(u(t)) = \int_{-\infty}^{\infty} \frac{u(\tau)}{\pi(t-\tau)} d\tau$$

linear operator  
Bendat & Piersol (1986)  
last chapter



How to get the Hilbert transform in practice

- (1) from ~~its~~ <sup>the</sup> Fourier coefficients (FFT) of the original data, i.e., get  $a(\omega), b(\omega)$
- (2) design a "filter" that changes phase by  $\pi/2$  but leaves amplitude intact



The complex covariance matrix

$$R_{pq} = \langle U_p^*(t) U_q(t) \rangle$$

represents an average of the cross-spectral matrix  $S_{pq}(f)$  over the entire frequency domain because

$$\begin{aligned}
R_{pq}(\tau=0) &= \int_{-\infty}^{\infty} R_{pq}(\tau) \delta(\tau=0) d\tau \\
&= \int_{-\infty}^{\infty} R_{pq}(\tau) \int_{-\infty}^{\infty} 1 \cdot e^{-j2\pi f\tau} df d\tau \\
&= \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} R_{pq}(\tau) e^{-j2\pi f\tau} df d\tau \\
&\stackrel{\text{handwriting}}{=} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} R_{pq}(\tau) e^{-j2\pi f\tau} d\tau df \\
&= \int_{-\infty}^{\infty} S_{pq}(f) df
\end{aligned}$$

The eigenvalue problem we now write again as

$$R_{pq} \cdot \vec{\phi} = \lambda \vec{\phi}$$

where we get (as in the standard EOF analysis) from the eigenvectors  $\vec{\phi}$

$$\begin{aligned}
U_p(t, x_p) &= \sum_{n=1}^M A_n(t) \cdot \phi_n^*(x_i) && \sum_{n=1}^M \phi_p(x_n) \phi_q^*(x_n) = \delta_{pq} \\
A_n(t) &= \sum_{p=1}^M U_p(x_p, t) \phi_n(x_i) && \frac{1}{K} \sum_{t=1}^K A_n(t_x) A_m^*(t_x) = \delta_{nm} \lambda_n
\end{aligned}$$

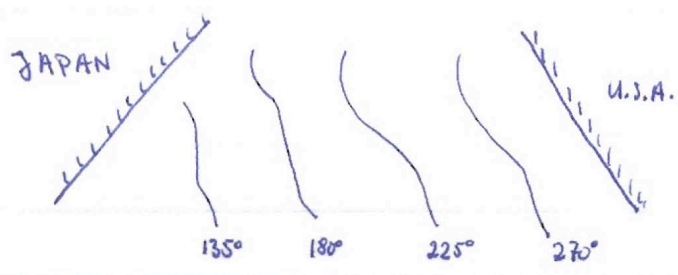
with

Because both eigenvectors and amplitude series are complex, we can now define/interpret those as 4 different functions such as

at each frequency  
for each eigenvector

1. a spatial phase

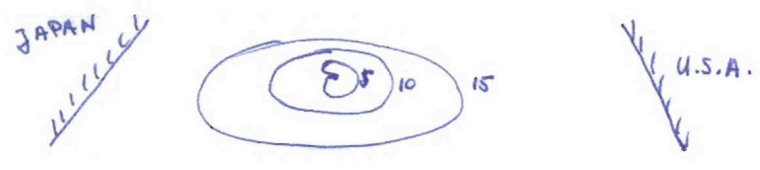
at each EOF mode n: 
$$\Theta_n(x_p) \equiv -\tan^{-1} \left[ \frac{\text{Im}(\phi_n(x_p))}{\text{Re}(\phi_n(x_p))} \right]$$



for a sinusoid  
 $u(x,t) = a \sin(kx - \omega_0 t)$   
there will be just 1 mode  
that goes through  $2\pi$   
over some distance  $L$   
i.e.  $\Theta_2 = 2\pi = \frac{2\pi}{L} x$

2. a spatial amplitude

$$B_n(x_p) = \sqrt{\phi_n(x_p) \phi_n^*(x_p)}$$

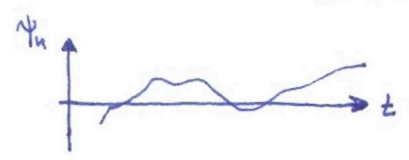


amplitude would  
be independent of x

3. a temporal phase

$$\Psi_n(t) = \tan^{-1} \left[ \frac{\text{Im}(A_n(t))}{\text{Re}(A_n(t))} \right]$$

also  
indicate  
non-stationarity  
to some extent



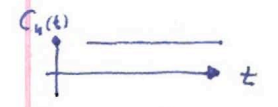
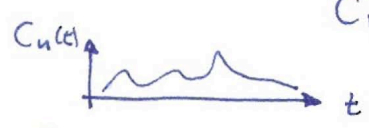
a pure sinusoid  
 $u(x,t) = a \sin(2\pi x - \omega_0 t)$   
would give  $\Psi_n(t)$  vs  $t$



} like  
complex  
dissipative  
?

4. a temporal amplitude

$$C_n(t) = \sqrt{A_n(t) A_n^*(t)}$$





# Summary EOF / FEOF / CEOF Analysis

Principal component analysis

Factor analysis

~~multiple~~

- economy of description                      data reduction; reveal redundancies
- orthogonalization                              transform data into a set of
- quantification of variance distribution      mutually uncorrelated data
- applies to non-stationary data
- separates spatial from temporal variability
- FEOF / CEOF detect propagating "pattern" also
- EOF                      detect standing "pattern" only
- FEOF works best if signal is concentrated in a narrow frequency band
- CEOF performance improves by band-pass filtering data
- interpretation of EOF / CEOF / FEOF not always easy  
     often need additional supporting data to  
     physically interpret the statistical decomposition
- active area of research; mostly meteorological application  
     (but then; oceanography always a few years/decades behind)