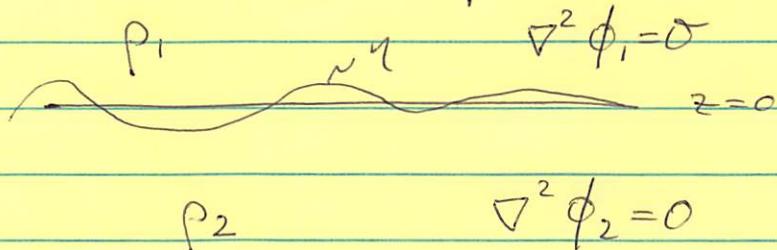


3.3 Internal Waves

Air - Atmosphere not the only interface
that can support gravity waves

Consider interface between two fluids of semi-infinite ~~ext~~ depth



At interface $z=0$ we have

$$\text{kinematic BC} \quad \eta_t = \phi_{1z} \quad \text{and} \quad \eta_t = \phi_{2z}$$

$$\begin{aligned} \text{dynamic BC} \quad & \rho_1 (\phi_{1t} + g_y) = \rho_2 (\phi_{2t} + g_y) \\ (\text{pressure matching}) \end{aligned}$$

These equations are satisfied with

$$\phi_1 = A_1 e^{-i\sigma t + ikx - kz}$$

$$\phi_2 = A_2 e^{-i\sigma t + ikx + kz}$$

$$\eta = a e^{-i\sigma t + ikx}$$

deep water
simplifies
 z -dependence

And the three surface conditions with these solutions become

$$-i\sigma a = -k A_1$$

$$-i\sigma a = +k A_2$$

$$\rho_1 \left(\frac{\sigma}{k} + g \right) = \rho_2 \left(-\frac{\sigma}{k} + g \right)$$

~~$\sigma^2 \approx g$~~

The pressure matching conditions can be written as

$$\sigma^2 = g k \frac{\rho_2 - \rho_1}{\rho_2 + \rho_1}$$

which for $\rho_1 = 0$

$$\sigma^2 = g k$$

which is the deep water dispersion relation

$$g \longrightarrow g \cdot \Delta p / \bar{\rho} \approx 10^{-3} g \quad \left| \begin{array}{l} \Delta p = \rho_2 - \rho_1 \\ \bar{\rho} = \rho_2 + \rho_1 \end{array} \right.$$

Much more on internal waves later

Let's look more closely at the nonlinear

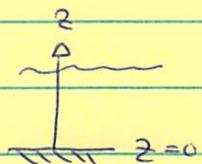
governing equations and boundary conditions

$$(1) \quad \phi_t + g\gamma + \frac{1}{2}(\phi_x^2 + \phi_y^2 + \phi_z^2) = 0 \quad @ z = D + \gamma$$

$$(2) \quad \gamma_t + \phi_x \gamma_x + \phi_y \gamma_y = \phi_z \quad @ z = D + \gamma$$

$$(3) \quad \nabla^2 \phi = 0$$

$$(4) \quad \phi_z = 0 \quad @ z = 0 \quad \text{which is the bottom}$$



Introduce non-dimensional variables

$$(x, y) = (x, y) \cdot L$$

$\underline{\Phi}$

$$z = z \cdot \underline{D}$$

$$t = t \cdot L / \sqrt{g D}$$

$$\gamma = \gamma \cdot a$$

$$\phi = \phi \cdot g a L / \sqrt{g D}$$

Then dimensional

$$(2) \quad \gamma_t + \phi_x \gamma_x = \phi_z$$

becomes

$$\frac{a \sqrt{g D}}{L} \gamma_t + \frac{g a L \cdot a}{L^2} \phi_x \gamma_x = \frac{g a L}{D \sqrt{g D}} \phi_z$$

or

$$\eta_t + \underbrace{\left(\frac{a}{D}\right)}_{\substack{\text{surface amplitude} \\ \text{depth}}} \phi_x \gamma_x = \underbrace{\left(\frac{L}{D}\right)^2}_{\delta} \phi_z @ z = 1 + \underbrace{\left(\frac{a}{D}\right)}_{\text{non-dimensional}} z$$

$\equiv \varepsilon \quad \equiv \delta$

parameters

With these, the other equations write as

$$(1) \phi_t + \frac{\varepsilon}{2} (\phi_x^2 + \phi_y^2) + \frac{\varepsilon}{2\delta^2} \phi_z^2 + \gamma = 0 @ z = 1 + \varepsilon y$$

$$(2) \eta_t + \varepsilon (\gamma_x \phi_x + \phi_y \gamma_y) = \frac{1}{\delta^2} \phi_z @ z = 1 + \varepsilon y$$

$$(3) \phi_{zz} + \delta^2 (\phi_{xx} + \phi_{yy}) = 0$$

$$(4) \phi_z = 0 @ z = 0$$

Case A : $\varepsilon \ll 1$ and $\delta = 1$

Liner Deep Water Waves

$$\left\{ \begin{array}{l} (1)+(2) \quad \phi_t + \gamma = 0 \quad \eta_t = \phi_z @ z = 1 \\ (3) \quad \nabla^2 \phi = 0 \\ (4) \quad \phi_z = 0 @ z = 0 \end{array} \right.$$